

Previous

More normal coords:

$$R_{ab}^c d = 2a^T_b{}^c d - 2b^T_a{}^c d$$

$$T_b^c d = \frac{1}{2} g^{ce} (g_{e,b,d} + g_{e,d,b} - g_{d,b,e})$$

$$R_{abcd} = \frac{1}{2} (g_{ad,bc} + g_{bc,ad} - g_{bd,ac} - g_{cd,ba} + g_{ac,bd} - g_{ad,bc})$$

Def'n w/ reps of \mathbb{Z}_2 in 2d rep

$$g_{\dots} \in \text{Fix}(12)$$

Extra rel'n (torsion-free)

$$(1 + (123) + (123)^2) g_{\dots} = 0 \quad \text{or } \sum_{\sigma \in \mathbb{Z}_3} g_{\dots}$$

From geod eq'n $2r = g_{red} r$

- 1-(12) ~~act~~
- 1-(34) ~~triv~~
- 1- $\sigma_1 \sigma_2$ (12)(34) ~~stid~~

$\Rightarrow g_{\dots} \in \text{Fix} \text{ also}$

$\Rightarrow (13)(24)$ acts trivially

$$\Rightarrow R_{abcd} = g_{ad,bc} - g_{ac,bd}$$

$$R = ((24) - (23)) g$$

Gauss's lemma: $r d_r = \nabla r$ $r d_r = \delta_{ij}^i x^j dx^i$

$\langle x^i d_i, \nabla r \rangle = g^{ik} \delta_{ij}^k x^j d_k r$

$x^i g_{ie} = \delta_{ie} x^i$

$x^i g_{ie} = x^i \delta_{ie} + \underbrace{x^i a_{ijk} x^k}_{=0}$

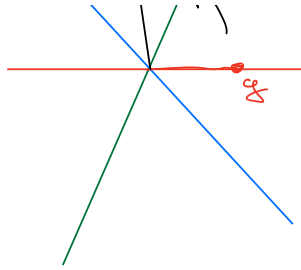
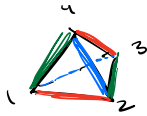
$\Leftrightarrow \sum_{e \neq i} a_{iie} = 0$ $d_r = x^i d_i$

$\frac{x^i d_i}{r} = \frac{g^{ij} x^j d_j}{r}$

$x^i = g^{ij} x^j$

$= x^i - g_{ij}^k x^k x^j x^i$





Normal coordinate trick: 2nd Bianchi identity

Use normal coords when you can to drop terms!

$$\nabla_a \nabla_b \nabla_c \omega_d - \nabla_b \nabla_a \nabla_c \omega_d = R(a,b)(\nabla \omega_d) = 0 \quad !$$

Recall the definition of \mathcal{R}

Con \mathcal{F} sec = C then

$$\begin{array}{r} \cancel{2} 3 1 - 1 3 2 \\ 1 2 3 - 2 1 3 \\ 3 1 2 - \cancel{3} 2 1 \end{array}$$

$$\mathcal{R}(x, y)z = C(\langle v, z \rangle x - \langle x, z \rangle y)$$

Check Ricci

$$\langle w, \mathcal{R}(x, y)z \rangle$$

(...)

$$((1-\sigma^2)(1-\langle z, z \rangle) = (1-\sigma^2)(1-\langle z, z \rangle) = 0$$

$$\text{let } u(x) = \begin{cases} t & C=0 \\ R \sin \frac{t}{R} & C = \frac{1}{R^2} > 0 \\ R \sinh \frac{t}{R} & C = -\frac{1}{R^2} < 0 \end{cases}$$

Then $(F(u, y))$ has const. sec. curvature C then it is locally isometric to $(\mathbb{R}^2, u_c(r)^2 g_{\text{Euc}})$

↑

compute $\mathcal{L}_x B$

$$B = \frac{1}{2g} \mathcal{L}_x g$$

$$\mathcal{L}_x B + B^2 = -R_x = \mathcal{R}(x, x) = -CI \quad \text{on } \mathbb{R}^2$$

$$\boxed{R_x = -\mathcal{R}(x, x) = \mathcal{R}(x, x)}$$

Hyperbolic plane models

Waveparticle

Note geodesics are easier to find than distance equations.
so phrasing things in terms of Jacobi
fields is a little more general.

negative energy
of geodesics
must stabilize!

Prop J is tangent to a variation thru geodesics iff

$$\nabla_t \nabla_t J = -R_{\cdot} J$$

Def'n to the Riccati eq'n

Γ picks out a Lagrangian subspace of the 2n-dim
space of Jacobi fields, namely $\{J \text{ s.t. } L_{\alpha} J = 0\}$

Remark $L_{\alpha} J = 0$ makes sense

$$\mathcal{L}(g^{\circ} \mathbb{R}) = \text{Der}(\mathcal{L}(\mathbb{R}), \mathcal{L}(\mathbb{R})_{\alpha})$$

$\partial_t J - \partial_t^2 F$ makes sense b/c $\partial_t \in \text{Der}(\mathcal{L}(\mathbb{R}))$ also

$$L_{\alpha} J = \nabla_r J - \nabla_S^2 F = \nabla_r J - B(S)$$

If $\nabla_r J = B(S)$ i.e. $L_{\alpha} J = 0$, then

$$(L_{\alpha} B)(S) + B^2 S = \nabla_r \nabla_r J$$

2nd von Borchers + group form. (7.1)

Sol'n's w/ $J(0) = 0$ determined by $\nabla_r J(0) \in T_{\alpha}$

\rightarrow unique on \mathbb{R}^n